Thermodynamics

2018/2019, lecturer: Martin Zápotocký

2 lectures:

- Thermodynamic processes, heat and work, calorimetry, 1st and 2nd law of thermodynamics
- 2. Entropy, thermodynamic potentials, nonequilibrium steady states, heat transport, heat balance in human organism

Warning: figures and equations drawn on the blackboard are not included in these slides

Thermodynamic system

- Thermodynamics studies processes of energy transformation on the macroscopic level (vs. kinetic theory on the molecular level)
- Thermodynamic system (containing a macroscopic number of molecules) is in contact with the surroundings (environment, reservoir). The system can be
 - isolated if no energy or matter is exchanged (e.g. liquid in perfectly insulating container)
 - closed if only energy is exchanged (e.g. gas constrained by piston)
 - in general open (e.g. fluid in contact with its vapor; living cell)
- State variables: pressure, temperature, volume, composition, ... uniquely determine the macroscopic state of the system

Thermodynamic process

- Thermodynamic process: a change in the state of the system
- Example: ideal gas constrained by movable piston can change p, V, T, constrained by equation of state pV = nRT
- Isobaric process: fixed *p*; isochoric: fixed *V*; isothermal: fixed *T*
- In isochoric process, no mechanical work performed
- In adiabatic process, no heat is exchanged with environment
- Reversible process: by reversing the process, one can bring both the system and its environment back into the intial state
- Whenever energy is dissipated (mechanical into heat, e.g. due to viscosity), the process is irreversible

Thermodynamic equilibrium

- A system in thermodynamic equilibrium would remain in the same state if we isolate it. There are no net fluxes.
- System is in thermodynamic equilibrium with environment if it has same T (thermal eq.), same p (mechanical eq.), same chemical potential (chemical eq.), and there are no unbalanced forces (fluxes).
- After a change in external conditions, the system requires a finite relaxation time to achieve the new state of thermodynamic equilibrium
- A reversible process goes through a sequence of equilibrium states; must be sufficiently slow (quasi-static)
- Living systems: maintained far from thermodynamic equilibrium (due to non-zero fluxes of matter and energy between system and environment)

Isothermal vs. adiabatic process, thermal reservoir

- Isothermal vs. adiabatic process:
 - example of isothermal process: system in contact with reservoir with large heat capacity + sufficiently slow process to maintain thermal equilibrium between system and reservoir
 - a sufficiently fast process will be adiabatic not enough time to exchange any heat with environment
- For ideal gas:
 - during isothermal process, pV = const
 - during adiabatic reversible process, $pV^{\gamma} = const$, $\gamma = \frac{c_p}{c_r} = 5/3$
- Newton's mistake when calculating the speed of sound

Heat vs. work

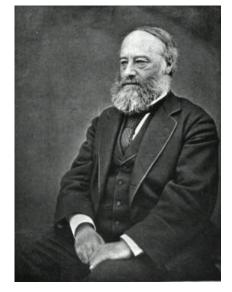
- Heat is the transfer of energy between two systems, during which no work is performed
- Spontaneous heat transfer from system with higher temperature T₁ to system with lower temperature T₂ (if the systems are in thermal contact).

 $Q_1 < 0$, $Q_2 > 0$ (system 2 gained heat)

- Basic unit for both heat and work is 1Joule = Newton × m = kg m²/s²
- Reminder:

If under constant force \vec{F} a body is moved by space interval \vec{x} , work done is

 $W = \vec{F} \cdot \vec{x} = F x \cos \alpha$



James Prescott Joule 1818-1889

Work expressed as integral

• For gas with piston of area A displaced by small interval dx: work done by the gas is

$$dW = \vec{F} \cdot d\vec{x} = p A dx = p dV$$

where dV = A dx is the volume change

• Therefore work done by gas during a process that changes the volume from V_i to V_f is given by the integral

$$W = \int_{V_i}^{V_f} p \, dV$$

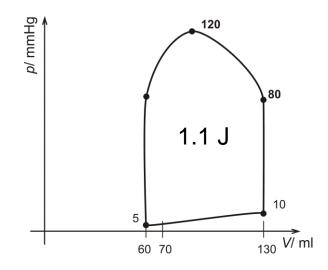
where the pressure p(V) can depend on volume V!

- Meaning of the integral: area under the p(V) curve
- This implies that W depends on the path of the process, i.e. work is not a state variable. E.g. for isothermal process in ideal gas,

$$W = nRT \ln(\frac{V_f}{V_i})$$

Example: work loop of the left ventricle

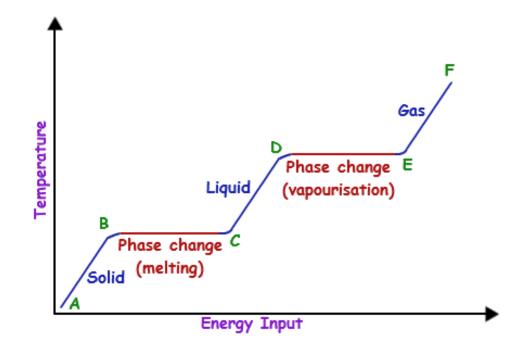
• The amount of work performed by the left ventricle during one cycle of heart contraction: the area within the loop in the diagram of intraventricular blood pressure vs. ventricle volume



- 1 contraction (in resting physical state) performs work of 1.1 J.
- Total power of the heart muscle (work/second): 1.3 W to 8 W, depending on level of physical activity. Regulated by heart rate and stroke volume.

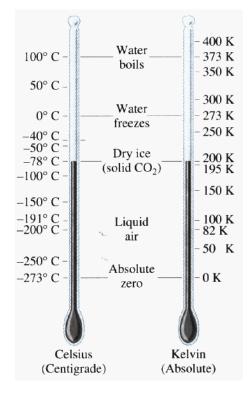
Heat is distinct from temperature!

- Temperature quantifies the kinetic energy of microscopic thermal motion (recall the equipartition theorem)
- In a given thermodynamics state, the system has a particular temperature T, but does not have a particular heat Q
- Heat is a transfer of energy during a process, and is not a state quantity



Measurement of temperature

- To measure temperature T, we bring the thermometer into thermal equilibrium with the measured system
- Temperature scales: defined by the zero point and the unit
- Celsius (centigrade) scale: zero point is the temperature of water (at standard p), unit = 1 °C, water boils at 100 °C.
- Fahrenheit scale (in ${}^{o}F$): $32 + \frac{9}{5}t_{Celsius}$
- Absolute scale: zero point is the lowest possible temperature (all thermal motion stops), absolute zero = 0 K = -273.15 °C. Unit is 1 Kelvin = 1 °C.
- Therefore the temperature of human body on the absolute scale: 273.15 °C + 37 °C = 310.15 K



Types of thermometers

- Liquid thermometers:
 - based on liquid volume expansion, $\Delta V/V_0 = \beta \Delta T$
 - the coefficient of expansion for mercury: $\beta = 1.8 \cdot 10^{-4} \text{ K}^{-1}$ (vacuum above liquid)
 - mercury thermometer range: -39 °C to +250 °C
- Bimetallic strip: difference in volume expansion coefficients of two metals welded together → bending (used in mechanical thermostats)
- Based on changes in electric resistance:
 - metal resistance thermometer (e.g. platinum, wide range)
 - thermistor (lab No. 3) strong decrease of resistance with T in a semiconductor (mK precision)
- Thermocouple: two metals in contact generate electromotive force (voltage) with magnitude depending on T. Measure voltage difference between hot and cold junction.
- Thermography: record intensity of thermal radiation with infrared camera. Helps to localize tumors.

Specific heat, latent heat

• The specific heat capacity gives the amount of heat needed to increase temperature by 1 Kelvin (while remaining in the same phase):

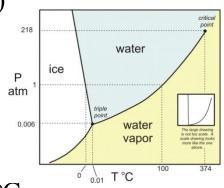
 $Q = mc\Delta T = nC\Delta T$

water: specific heat $c = 4.2 \text{ kJ kg}^{-1}\text{K}^{-1}$, molar specific heat C = Mc

- Specific heat at constant p is higher than at constant V:
 - due to work $p\Delta V$
 - for ideal gas: $C_P = C_V + R = \frac{3}{2}R + R = \frac{5}{2}R$
 - adiabatic index $\gamma = C_P/C_V > 1$ (water vapors ≈ 1.3)
 - for fluids close to 1
- Latent heat is the heat associated with a phase transition

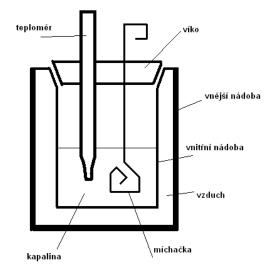
- Q = mL, L = specific latent heat,

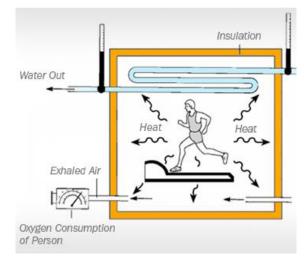
- e.g. $L = 2.3 \text{ MJ kg}^{-1}$ for evaporation of water at 100 °C



Calorimetry

- Heat can be measured in a calorimeter. Based on measurement of temperature change + knowledge of heat capacity.
- Basic calorimeter: water in thermally isolating vessel
- If water temperature increases by ΔT :
 - water absorbed heat $Mc\Delta T$ (M = mass of water, c = specific heat of water)
 - rest of calorimeter absorbed heat $Kc\Delta T$, where K = water value of the calorimeter (kg) calibrated before measurement
 - total heat absorbed (calorimetric equation): $Q = (M + K) c \Delta T$
- Calorimetric units: 1 cal = heat needed to bring 1 g of water from 14.5 °C to 15.5 °C; 1 cal = 4,187 J
- Direct calorimetry of metabolic rate whole body calorimeter





Internal energy

• Internal energy of a system:

internal kinetic energy + potential energy of internal forces

- Internal kinetic energy: sum of kinetic energies of all system parts
 - does not include motion of the whole system (center of mass)
 - includes translational, rotational, and vibrational kinetic energy of molecules
 - average kinetic energy of 1 molecule = $n \frac{1}{2k_B}T$; v ~ 1 km/s

(n = number of degrees of freedom)

 $k_{\rm B}$ = Boltzmann const = 1.38×10⁻²³ J/K)

- Internal potential energy includes the energy of chemical bonds; does not include energy from external forces (e.g. gravitation)
- Example: internal energy of 1 mole of inert gas (no chemical bonding): $N_A 3 \frac{1}{2} k_B T = 3/2 RT = 3406 J$

First law of thermodynamics

- Law of conservation of energy: the total energy of a isolated system does not change in time. Note: the only truly isolated system is the Universe!
- A thermodynamic system in interaction with environment: exchange of energy in the form of work and heat
- 1st law of thermodynamics:

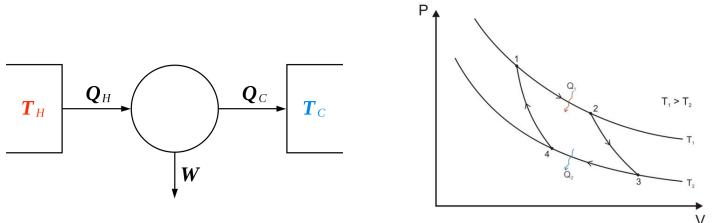
 $\Delta U = Q - W$

where $\Delta U = U_f - U_i$ is the change in internal energy, Q the heat received and W the work performed during the process

- Internal energy is a state variable! ΔU does not depend on the path of the process, even though Q and W are not state variables.
- Remember the signs: received heat increases internal energy ($\Delta U > 0$), performed work decreases internal energy ($\Delta U < 0$). For an isolated system, $\Delta U = 0$.

Cyclic process, heat engine

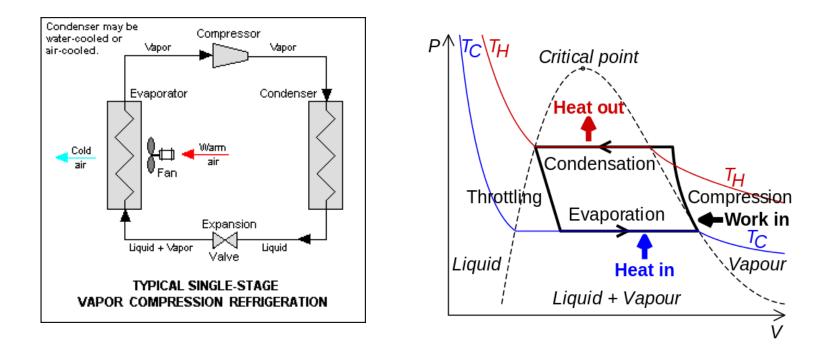
- A cyclic process returns to the same state
- Work in cycle is equal to the area enclosed by the process path in the p-V diagram (positive if clockwise)
- A heat engine performs conversion of heat energy to mechanical work. Two segments of the cycle (1-2 and 3-4) operate at temperatures of hot and cold reservoir



• E.g. Carnot cycle: isothermal expansion (heat intake) → adiabatic expansion (work output) → isothermal compression → adiabatic compression

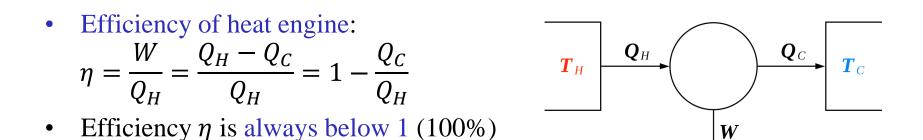
Heat engine running in reverse: the heat pump

• In a vapor compression refrigerator: the coolant absorbs heat when it is evaporating during expansion



Second law of thermodynamics

• It is impossible to make a heat engine which only takes heat from one reservoir. Part of the heat must be passed to another reservoir with lower temperature. (It is therefore degraded – cannot be put back into hot reservoir without additional energy cost.)



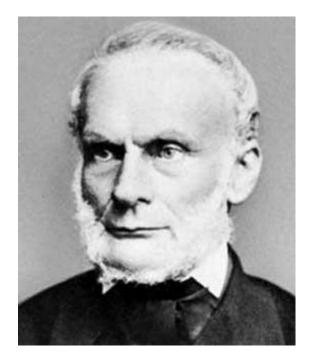
• The maximum possible efficiency of a heat engine is

$$\eta = 1 - \frac{I_C}{T_H}$$

and it is reached only if the engine process is reversible

• Alternative formulation of 2nd law of thermodynamics: using entropy.

History of entropy: macroscopic and microscopic



<text>

Rudolf Clausius,1865: The energy of the universe is constant. The entropy of the universe tends to a maximum. Ludwig Boltzmann,1877: entropy measures disorder $S = k_B \ln W$

Entropy

• Entropy expresses the degree of disorder in the system. Given by: $S = k_B \ln W$

 $k_{\rm B}$ = Boltzmann const = 1.38×10⁻²³ J/K

W = thermodynamic likelihood of the state

(number of microscopic ways to obtain the macroscopic state)

- Increase in *S* means that the new state of the system is less organized.
- Example: melting of the crystal lattice increases entropy.
- Example: 1 mol of ideal gas put in left half of partitioned container. After removal of partition, the gas freely expands into right half, and $\Delta S = k_B \ln \frac{W_f}{W_i} = k_B \ln 2^{N_A} = k_B N_A \ln 2 = R \ln 2 = 5,76 \text{J/K}$
- 2nd law formulated in terms of entropy: In an isolated system: entropy can never decrease. It will increase during any irreversible process. The system will spontaneously approach the most probable configuration: the state of thermodynamic equilibrium, with maximal entropy.

Heat and entropy

• During a reversible isothermal process, the entropy change is proportional to absorbed heat:

$$\Delta S = S_f - S_i = \frac{Q}{T}$$

• For any irreversible process:

$$S_f - S_i > \int_i^f \frac{dQ}{T}$$

therefore for any irreversible adiabatic process, $\Delta S > 0$

• Therefore the entropy of an isolated system (the Universe) stays constant or increases. However, in a non-isolated subsystem, entropy can decrease (at the expense of the rest of the system).

Dissipation, "lost work" in irreversible process

• For reversible process at temperature T:

$$\Delta U = Q_{\rm rev} - W_{\rm rev} = T\Delta S - W_{\rm rev}$$

- Consider an irreversible process with the same initial and final states: $Q_{irr} - W_{irr} = \Delta U = T\Delta S - W_{rev}$ Since $Q_{irr} < T\Delta S$, we have $W_{irr} < W_{rev}$
- Work done in the irreversible process is lower by the ,,lost work" amount $W_{lost} = W_{rev} - W_{irr}$. Interpretation: some of the internal energy is dissipated, i.e. used to increase entropy rather than to perform work.
- Can write for any process at temperature T:

$$\Delta S = \frac{Q}{T} + \frac{W_{\text{lost}}}{T}$$

i.e. entropy of the system can be changed by heat transfer with surroundings (*Q*) or by internal dissipation ($W_{lost} > 0$)

Thermodynamic potentials

- The 4 thermodynamic potentials express the potential of the system to do work under specific conditions
- Are minimized in thermodynamic equilibrium
- Analogous concept in mechanics: potential energy
- State functions, do not depend on process path. Extensive quantities, units: Joule.
- Internal energy U
 - recall: internal kinetic energy + potential energy of internal forces
 - *U* = energy needed to create the system (in isolation)
 - $\Delta U = Q W$, therefore in an adiabatic process, Q = 0, $W = \Delta U$
 - for isolated system at fixed volume, U is minimized in thermodynamic equilibrium

Enthalpy, free energy, free enthalpy

- Free energy (Helmholtz function) F = U TS
 - energy needed to create the system while in contact with reservoir at T
 - for isothermal, reversible process: $\Delta F = \Delta U T\Delta S = \Delta U Q = -W$
 - i.e. decrease in free energy = maximum work that can be done in a process at fixed T and V. Less work done if the process is irreversible.
- Enthalpy H = U + pV
 - for isobaric process: $\Delta H = \Delta U + p\Delta V = \Delta U + W = Q$,
 - i.e. change in enthalpy = heat received by system
 - for exothermal reaction at constant pressure, Q < 0, enthalpy decreases (and achieves minimum in thermodynamic equilibrium)
 - Law of Hess: reaction heat does not depend on the reaction path (enthalpy is a state function)
- Free enthalpy (Gibbs function) G = U + pV TS = H TS
 - if $\Delta G < 0$, reaction will be spontaneous (at const T and p)
 - G is the maximum work that can be performed by chemical reaction

Chemical potential

(is not a thermodynamic potential – not minimized in TD equilibrium)

- In an open system, changes in composition lead to changes in energy
- Chemical potential μ_i for chemical substance i: change in energy of the system if 1 mole of substance i is added (and the amount of other substances is not changed)
- At given p, T: change in free enthalpy of the equilibrium state is $\Delta G = \mu_1 \Delta n_1 + \mu_2 \Delta n_2 + \cdots$

where Δn_1 , Δn_2 , ... are the changes in number of moles

• Chemical potential is a measure of chemical affinity, but depends on state parameters. For ideal gas

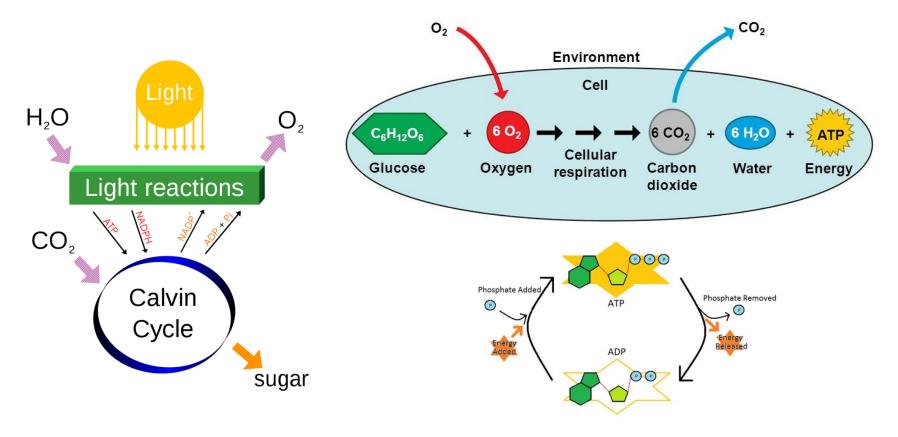
$$\mu(p,T) = \mu_o(T) + k_B T \ln \frac{p}{p_0}$$

Nonequilibrium steady states, principle of minimal entropy production

- Living systems operate far from thermodynamic equilibrium and continously dissipate energy
- ΔS (of system + environment) per unit time = entropy production rate > 0
- Nonequilibrium steady state examples:
 - steady heat flux in system in contact with two thermal reservoirs
 - combustion engine: steady supply of reactants + steady removal of products
- The thermodynamic potentials are not minimized in such noneq. states. Under restricted conditions (not too far from equilibrium), the following principle holds:
- Principle of minimal entropy production: a nonequilibrium system tends to lower its rate σ of production of entropy (per unit volume), and reaches a minimum σ in the nonequilibrium steady state.

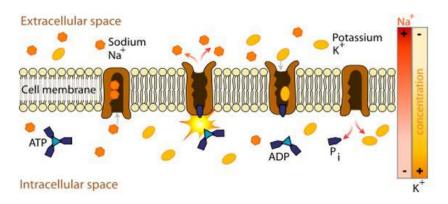
Energy conversions in biological processes

- Remember from biochemistry:
 - Photosynthesis in chloroplasts
 - Cellular respiration: nutrients \rightarrow ATP
 - ATP hydrolysis: intracellular source of energy (7.3 kcal/mol)



Passive and active transport at the cell membrane

- Passive transport:
 - down the concentration gradient, or more generally down the electrochemical gradient (see lectures on electric phenomena)
 - requires no energy (energy is gained during the transport)
 - examples: diffusion of O_2 into cell; facilitated diffusion of glucose from blood into eryhtrocytes; ions flowing through an open ion channel
- Active transport:
 - against diffusion / electrical force. Consumes energy supplied by ATP hydrolysis.
 - Example: in each cycle, the Na⁺/K⁺ pump moves 3 Na⁺ out and 2 K⁺ in (against concentration gradient), hydrolysing 1 ATP



Energetic requirements of human organism

- Most energy is obtained through metabolic reactions (food). ATP used in nonequilibrium (active) processes inside cells. See notes p. 58-60.
- Metabolic energy usage, per unit area of body surface (total area 1,7 m²)
 - sleep: $35 \text{ kcal/m}^2/\text{hod}$
 - reading, sitting: $60 \text{ kcal/m}^2/\text{hod}$
 - walking 5 km/h: 140 kcal/m²/hod
 - cycling
 250 kcal/m²/hod
 - running $600 \text{ kcal/m}^2/\text{hod}$
- For an active average human, total of 2300 kcal/day = 11 MJ /day
- Most of the metabolic energy is converted to heat (e.g. during muscle contraction, less than 40% of energy from ATP hydrolysis is converted to mechanical work).

Mechanisms of heat transport

• Heat conduction: thermal energy of molecular motion is transferred by molecular collisions. The rate of heat transfer (amount of heat transferred per unit time) is proportional to temperature gradient:

$$\frac{dQ}{dt} = -\kappa A \frac{dT}{dx}$$

A = area of surface perpendicular to \hat{x} , κ = thermal conductivity

- For water: $\kappa = 0.6 \text{ Wm}^{-1}\text{K}^{-1}$, air: 0.025 Wm⁻¹K⁻¹, metal: 100 Wm⁻¹K⁻¹ (high κ results in cold feeling)
- Convection: occurs in fluids. Decreased density in locations with lower T
 → buoyancy forces move fluid, bulk flow transports heat
- Radiation: vibrating charged particles emit electromagnetic radiation (act as antenna). Internal energy of the system is transformed into energy of radiation. Radiated energy per unit time, from surface of area A:

$$\frac{dQ}{dt} = e\sigma AT^4$$

 σ = Stefan-Boltzmann constant = 5.67 · 10⁻⁸ Wm⁻²K⁻⁴

e = emissivity of the material ($e \approx 1$ for a black body = perfect absorber)

Thermal balance in human organism

- More than 80% of energy is transformed into heat. Efficient mechanisms for heat removal required to prevent overheating and death.
- Most heat is generated deep in tissues. The coefficient of heat conduction in a tissue is far too low to permit conduction of heat to body surface (would conduct only 20 kcal/hour)
- More efficient mechanism: heat conduction from cells do blood capillaries, then heat convection to body surface
- The skin loses heat:
 - through convection in surrounding air. Depending on air velocity, 65 kcal/hour to several hundred kcal/hour.
 - through radiation. The net heat loss = difference between energy radiated and absorbed = $Ae\sigma(T^4 T_{ok}^4)$, with skin emissivity e ≈ 1 (maximum radiation at 10 µm wavelength, i.e. infrared). At 7 °C temperature difference between skin and surroundings, the rate of radiative heat loss is 60 kcal/hour.

Heat loss through evaporation

- This is still not enough... In warm weather (above 31 °C), or even cold weather while excercising, an additional mechanism of heat loss becomes necessary to prevent overheating: evaporation of sweat.
- Latent heat for sweat vaporization (at 37 °C) is 580 kcal/kg.
- The human organism can produce up to 4 litres of sweat within 1 hour, and can keep producing 1 litre/hour (assuming rehydration)
- Only the sweat that evaporates leads to efficient heat loss! Evaporated portion is highest in hot, dry, windy weather. Cannot vaporize if relative air humidity close to 100% (saturated vapor).